

## New Disordering Mode for TFSI<sup>-</sup> Anions: The Nonequilibrium, Plastic Crystalline Structure of Et<sub>4</sub>NTFSI

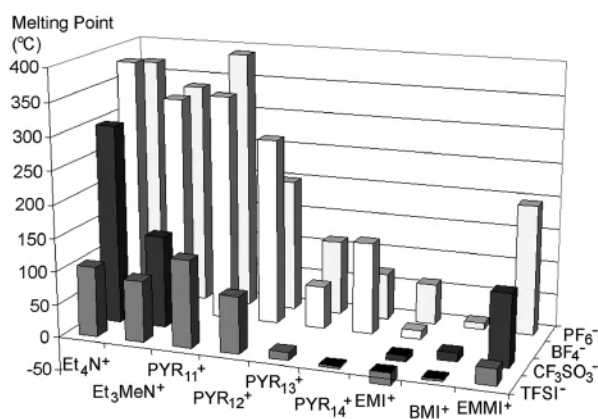
Wesley A. Henderson,<sup>†</sup> Marie Herstedt,<sup>‡</sup> Victor G. Young, Jr.,<sup>§</sup> Stefano Passerini,<sup>||</sup> Hugh C. De Long,<sup>⊥</sup> and Paul C. Trulove<sup>\*†</sup>

Department of Chemistry, U.S. Naval Academy, Annapolis, Maryland 21402, Laboratoire de Physico-Chimie Moléculaire, UMR 5803, CNRS, Université Bordeaux I, Talence Cedex 33405, France, X-ray Crystallographic Laboratory, Department of Chemistry, University of Minnesota, Minneapolis, Minnesota 55455, Casaccia Research Center, ENEA, Rome, Italy 00060, and Air Force Office of Scientific Research, Arlington, Virginia 22203

Received August 11, 2005

A new TFSI<sup>-</sup> anion disordering mode has been discovered in a supercooled plastic crystalline phase of Et<sub>4</sub>NTFSI, which may, in part, account for the low melting points of TFSI<sup>-</sup> salts with organic cations, thereby forming ionic liquids, and the intriguing properties of LiTFSI for lithium battery applications.

The factors that determine the melting point of a salt, and thus the formation of room-temperature ionic liquids, remain poorly understood. Figure 1 displays the melting points of salts with commonly utilized cations and varying anions.<sup>1</sup> It is intriguing that the N(SO<sub>2</sub>CF<sub>3</sub>)<sub>2</sub><sup>-</sup> (TFSI<sup>-</sup>) anion is so effective at reducing a salt's melting point, whereas the same effect is not found for other charge-delocalized anions. This has been attributed to the lower symmetry and bulky nature of the anion (relative to, for example, BF<sub>4</sub><sup>-</sup>), as well as its extensive charge delocalization<sup>2</sup> and flexibility.<sup>3</sup> In particular, two low-energy conformations are known with C<sub>2</sub> and C<sub>1</sub>



**Figure 1.** Melting points of various salts (PYR<sub>13</sub><sup>+</sup> = *N*-methyl-*N*-propylpyrrolidinium, EMI<sup>+</sup> = 1-ethyl-3-methylimidazolium, BMI<sup>+</sup> = 1-butyl-3-methylimidazolium, and EMMI<sup>+</sup> = 1-ethyl-2,3-dimethylimidazolium).

symmetry (Figure 2a). Here we report a nonequilibrium, disordered structure for Et<sub>4</sub>NTFSI, which reveals a new disordering mode for the TFSI<sup>-</sup> anion.<sup>4</sup> This disordering mode may help explain both the “crystallinity gap” in poly-(ethylene oxide)–LiTFSI electrolytes<sup>5</sup> and the low melting points of organic salts with TFSI<sup>-</sup> anions (Figure 1).

Et<sub>4</sub>NTFSI exhibits two solid-phase transitions at 277 and 322 K prior to melting at 377 K (Figure 2b). There is a strong hysteresis in the phase III → II transition near 277 K, which enabled a single crystal in the room-temperature phase II configuration to be supercooled to 100 K and the structure

\* To whom correspondence should be addressed. E-mail: trulove@usna.edu.

<sup>†</sup> U.S. Naval Academy.

<sup>‡</sup> CNRS.

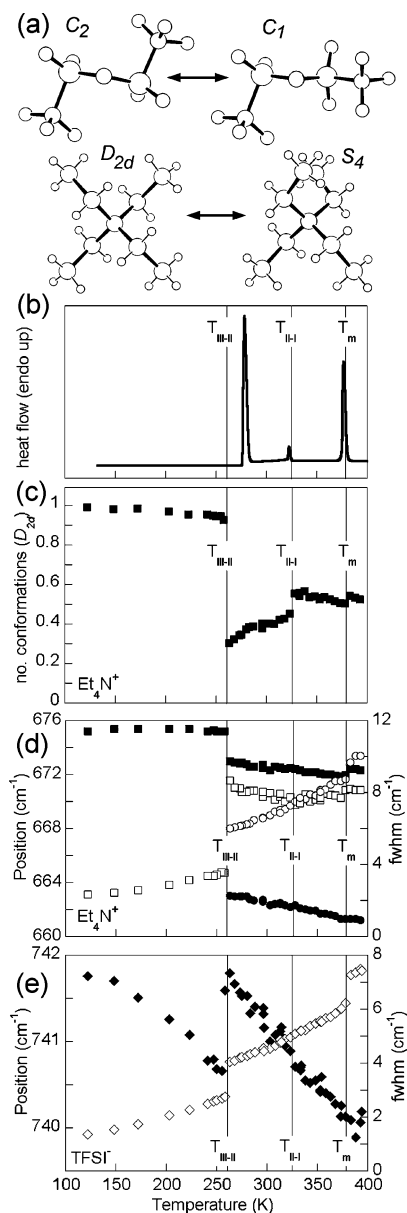
<sup>§</sup> University of Minnesota.

<sup>||</sup> ENEA.

<sup>⊥</sup> Air Force Office of Scientific Research.

- (1) (a) Xu, K.; Ding, M. S.; Jow, T. R. *J. Electrochem. Soc.* **2001**, *148*, A267. (b) Golding, J.; Hamid, N.; MacFarlane, D. R.; Forsyth, M.; Forsyth, C.; Collins, C.; Huang, J. *Chem. Mater.* **2001**, *13*, 558. (c) Forsyth, S.; Golding, J.; MacFarlane, D. R.; Forsyth, M. *Electrochim. Acta* **2001**, *46*, 1753. (d) MacFarlane, D. R.; Meakin, P.; Sun, J.; Amini, N.; Forsyth, M. *J. Phys. Chem. B* **1999**, *103*, 4164. (e) Henderson, W. A.; Passerini, S. *Chem. Mater.* **2004**, *16*, 2881. (f) Matsumoto, K.; Hagiwara, R.; Yoshida, R.; Ito, Y.; Mazej, Z.; Benkić, P.; Žemva, B.; Tamada, O.; Yoshino, H.; Matsubara, S. *Dalton. Trans.* **2004**, 144. (g) Ngo, H. L.; LeCompte, K.; Hargens, L.; McEwen, A. B. *Thermochim. Acta* **2000**, *357–358*, 97. (h) Bonhôte, P.; Dias, A.-P.; Armand, M.; Papageorgiou, N.; Kalyanasundaram, K.; Grätzel, M. *Inorg. Chem.* **1996**, *35*, 1168. (i) Huddleston, J. G.; Visser, A. E.; Reichert, W. M.; Willauer, H. D.; Broker, G. A.; Rogers, R. D. *Green Chem.* **2001**, *3*, 156. (j) Fredlake, C. P.; Crosthwaite, J. M.; Hert, D. G.; Aki, S. N. V. K.; Brennecke, J. F. *J. Chem. Eng. Data* **2004**, *49*, 954.
- (2) Benrabah, D.; Arnaud, R.; Sanchez, J.-Y. *Electrochim. Acta* **1995**, *40*, 2437.

- (3) (a) Johansson, P.; Gejji, S. P.; Tegenfeldt, J.; Lindgren, J. *Electrochim. Acta* **1998**, *43*, 1375. (b) Arnaud, R.; Benrabah, B.; Sanchez, J.-Y. *J. Phys. Chem.* **1996**, *100*, 10882. (c) Foropoulos, J., Jr.; DesMarteau, D. D. *Inorg. Chem.* **1984**, *23*, 3720.
- (4) There is one previous report of similar TFSI<sup>-</sup> anion disordering: Schlueter, J. A.; Geiser, U.; Wang, H. H.; Kini, A. M.; Ward, B. H.; Parakka, J. P.; Daugherty, R. G.; Kelly, M. E.; Nixon, P. G.; Winter, R. W.; Gard, G. L.; Montgomery, L. K.; Koo, H.-J.; Whangbo, M.-H. *J. Solid State Chem.* **2002**, *168*, 524.
- (5) (a) Vallée, A.; Besner, S.; Prud'homme, J. *Electrochim. Acta* **1992**, *37*, 1579. (b) Edman, L.; Ferry, A.; Doeff, M. M. *J. Mater. Res.* **2000**, *15*, 1950. (c) Tarascon, J.-M.; Armand, M. *Nature* **2001**, *414*, 359.



**Figure 2.** (a) Low-energy TFSI<sup>-</sup> anion and Et<sub>4</sub>N<sup>+</sup> cation conformations. (b) DSC heating trace of Et<sub>4</sub>NTFSI. (c) Fraction of Et<sub>4</sub>N<sup>+</sup> cation D<sub>2d</sub> conformers calculated from the 675 (D<sub>2d</sub>) and 663 (S<sub>4</sub>) cm<sup>-1</sup> Raman doublet. (d) Band position (filled symbols) and fwhm (open symbols) changes for the Et<sub>4</sub>N<sup>+</sup> cations (D<sub>2d</sub>, squares; S<sub>4</sub>, circles). (e) Band position (filled diamonds) and fwhm (open diamonds) changes for the TFSI<sup>-</sup> Raman band as a function of temperature.

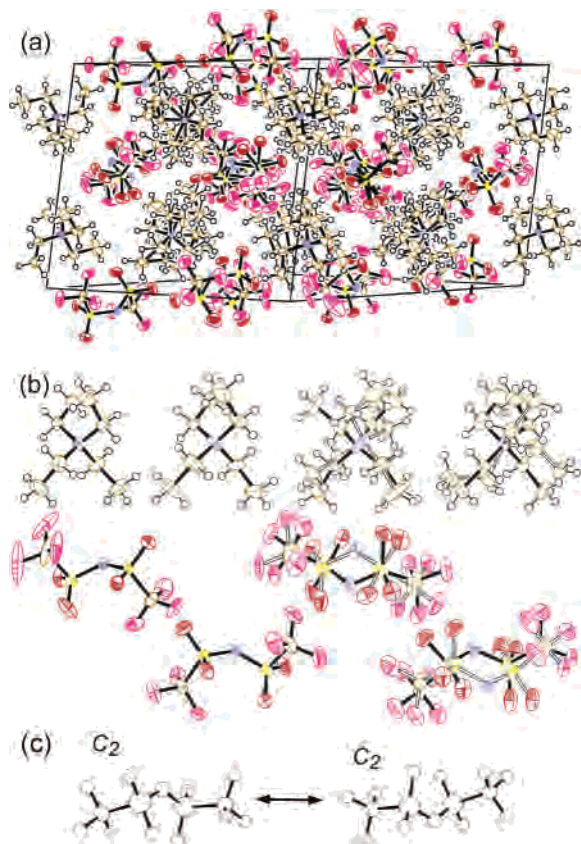
of this nonequilibrium phase determined. The hysteresis and shape of the differential scanning calorimetry (DSC) endothermic peak of Et<sub>4</sub>NTFSI at 277 K indicate that this is a first-order phase transition that results in a considerable structural change from the low-temperature phase III structure. The energy associated with this transition is greater than that of the fusion transition ( $\Delta_{\text{III-II}}H = 18.3 \text{ kJ mol}^{-1}$ ,  $\Delta_{\text{III-II}}S = 65.6 \text{ J K}^{-1} \text{ mol}^{-1}$ ;  $\Delta_{\text{I-fus}}H = 9.0 \text{ kJ mol}^{-1}$ ,  $\Delta_{\text{I-fus}}S = 24.0 \text{ J K}^{-1} \text{ mol}^{-1}$ ), suggesting that the ions become highly disordered at room temperature. Both the Et<sub>4</sub>N<sup>+</sup> cations and TFSI<sup>-</sup> anions may be examined by Raman spectroscopy to glean insight into the variations in ion conformations, disorder, and interactions with neighboring ions.<sup>6</sup> The hysteresis in the phase III  $\rightarrow$  II transition is evident in the

discrepancy between the DSC and Raman transition onsets (Figure 2).

The Et<sub>4</sub>N<sup>+</sup> cations are known to be flexible, adopting either an all-*trans* (tt,tt or D<sub>2d</sub>) Greek cross or *trans-gauche* (tg,-tg or S<sub>4</sub>) Nordic cross conformation in crystalline salts (Figure 2a).<sup>7,8</sup> A  $\nu_s(\text{C}_4\text{N})$  stretching Raman vibrational band near 675 cm<sup>-1</sup> originates from the D<sub>2d</sub> conformation, whereas the S<sub>4</sub> conformation gives a band near 663 cm<sup>-1</sup>.<sup>8</sup> Figure 2c shows the variation with temperature of the fraction of the Et<sub>4</sub>N<sup>+</sup> cations with the D<sub>2d</sub> conformation. At temperatures below the phase III  $\rightarrow$  II transition, nearly all of the Et<sub>4</sub>N<sup>+</sup> cations are ordered in the D<sub>2d</sub> conformation. After the transition (to the disordered phase II), however, only 30% of the cations adopt the D<sub>2d</sub> conformation (the remainder adopt the S<sub>4</sub> conformation). Figure 2d shows the variation with temperature of the  $\nu_s(\text{C}_4\text{N})$  stretching band positions and full-width at half-maximum (fwhm). The fwhm gives an indication of the distribution of ion conformations and thus of the disorder of the system. The position of the D<sub>2d</sub> band in the low-temperature phase III is constant with increasing temperature, but the fwhm increases somewhat. At the first phase transition (phase III  $\rightarrow$  II), the position of the D<sub>2d</sub> band decreases sharply and thereafter decreases slightly with increasing temperature. The fwhm increases sharply at this transition, indicating a large disordering of the cations. Thereafter, the fwhm of this band also decreases slightly with increasing temperature. The 663-cm<sup>-1</sup> band appears after the first phase transition, indicating the presence of cations with the S<sub>4</sub> conformation. The band position decreases steadily with temperature, but the fwhm increases strongly. The Et<sub>4</sub>N<sup>+</sup> cations are thus ordered in the D<sub>2d</sub> conformation in the low-temperature phase III, but become highly disordered in phase II.

As mentioned above, TFSI<sup>-</sup> anions are also highly flexible, with two low-energy transoid and cisoid conformations with C<sub>2</sub> and C<sub>1</sub> symmetry, respectively.<sup>3,9</sup> In this case, the Raman spectra aid in determining the degree of interaction between the TFSI<sup>-</sup> anions and neighboring cations. The TFSI<sup>-</sup> anion has a strong Raman band near 740 cm<sup>-1</sup> associated with the expansion and contraction of the entire anion.<sup>10</sup> This band is often used to identify coordination of the TFSI<sup>-</sup> anions with Li<sup>+</sup> cations.<sup>11</sup> Figure 2e shows the variation with temperature of the position and fwhm of this band. With increasing temperature, the band position decreases and the fwhm increases until the phase III  $\rightarrow$  II transition, at which point a dramatic upward shift occurs for both the anion band

- (6) Herstedt, M.; Henderson, W. A.; Smirnov, M.; Ducasse, L.; Servant, L.; Talaga, D.; Lassègues, J. C. *J. Raman Spectrosc.* **2006**, in press.
- (7) Brand, H. V.; Curtiss, L. A.; Iton, L. E.; Trouw, F. R.; Brun, T. O. *J. Phys. Chem.* **1994**, *98*, 1293.
- (8) Naudin, C.; Bonhomme, F.; Bruneel, J. L.; Ducasse, L.; Grondin, J.; Lassègues, J.-C.; Servant, L. *J. Raman Spectrosc.* **2000**, *31*, 979.
- (9) Holbrey, J. D.; Reichert, W. M.; Rogers, R. D. *Dalton Trans.* **2004**, 2267.
- (10) (a) Rey, I.; Johansson, P.; Lindgren, J.; Lassègues, J. C.; Grondin, J.; Servant, L. *J. Phys. Chem. A* **1998**, *102*, 3249. (b) Herstedt, M.; Smirnov, M.; Johansson, P.; Chami, M.; Grondin, J.; Servant, L.; Lassègues, J. C. *J. Raman Spectrosc.* **2005**, *36*, 762.
- (11) (a) Rey, I.; Lassègues, J. C.; Grondin, J.; Servant, L. *Electrochim. Acta* **1998**, *43*, 1505. (b) Brouillette, D.; Irish, D. E.; Taylor, N. J.; Perron, G.; Odziemkowski, M.; Desnoyers, J. E. *Phys. Chem. Chem. Phys.* **2002**, *4*, 6063. (c) Edman, L. *J. Phys. Chem. B* **2000**, *104*, 7254.



**Figure 3.** (a) Unit cell of the nonequilibrium, supercooled (phase II) crystal structure of Et<sub>4</sub>NTFSI at 100 K. (b) Views of the ions found in the asymmetric unit. Two of the cations and two of the anions are disordered (N, blue; S, yellow; O, red; F, pink). (c) The two C<sub>2</sub> anion conformations for the disordered TFSI<sup>-</sup> anions in Et<sub>4</sub>NTFSI.

position and fwhm. After the first transition, the wavenumber again decreases and the fwhm continues to increase until melting. This indicates that the TFSI<sup>-</sup> anions are also disordered in phase II. The variation in the 740-cm<sup>-1</sup> band and fwhm below the phase III → II transition suggests that, as the temperature increases, approaching the transition temperature, the anions become increasingly disordered and the interaction of the anions with the neighboring cations becomes weaker. This may, in fact, be responsible for inducing the phase transition.

A single crystal of the salt was cooled at a rate of 0.5 K min<sup>-1</sup> to 100 K, and the crystal structure was determined from the diffraction data. The structure at 100 K is both twinned and disordered (Figure 3a). The asymmetric unit consists of four Et<sub>4</sub>N<sup>+</sup> cations and four TFSI<sup>-</sup> anions (Figure 3b). Two of the cations and two of the anions are disordered. The two ordered cations adopt the S<sub>4</sub> conformation. The first disordered cation adopts D<sub>2d</sub> and S<sub>4</sub> conformations with 50.9% and 49.1% occupancy. The second disordered cation adopts the S<sub>4</sub> conformation with 83.0% occupancy and a second conformation of C<sub>1</sub> symmetry<sup>7</sup> with 17.0% occupancy. The structure is indeed the supercooled room-

temperature phase II structure because the fraction of Et<sub>4</sub>N<sup>+</sup> cations with the D<sub>2d</sub> conformation is only 12.7% (rather than close to 100% for phase III). The supercooling of the phase II salt into a nonequilibrium, frozen-in disordered phase is analogous to the freezing of a disordered liquid into a disordered glassy solid. The phase II equilibrium between the Et<sub>4</sub>N<sup>+</sup> cations with the D<sub>2d</sub> and S<sub>4</sub> conformations changes continuously with temperature. If the data for the fraction of cations with the D<sub>2d</sub> conformation (Figure 3c) are extrapolated from room temperature to about 200 K (near the expected glass transition), a value of about 13% is obtained. This indicates that, on cooling, the phase II structure is retained (becoming supercooled) and ionic motion is eventually completely frozen out, locking the ions into a particular conformation. All of the ions may be disordered at room temperature. The value of 200 K is close to the glass transition temperature of many TFSI<sup>-</sup> salts when they are quenched into supercooled amorphous phases.

Surprisingly, the anion disorder in the crystal structure does not consist of a mixture of the C<sub>2</sub> and C<sub>1</sub> conformations as expected.<sup>3,9</sup> Instead, all of the anions have a C<sub>2</sub> conformation, but two of the anions are disordered as two C<sub>2</sub> conformations with major occupancies of 60.8% and 77.5%, respectively, in which the nitrogen atom and SO<sub>2</sub>CF<sub>3</sub> groups are rotated but the SCF<sub>3</sub> groups retain their approximate positions (though rotated about the C<sub>3</sub> axis) (Figure 3b). This indicates a new dynamic mode for the disordering of the TFSI<sup>-</sup> anions not previously modeled or observed experimentally in which the anions flex within the cage created by the neighboring cations (Figure 3c). The same disordering mode has also been observed by us in crystal structures of PYR<sub>12</sub>TFSI,<sup>12</sup> PIP<sub>12</sub>-TFSI (with *N*-ethyl-*N*-methylpyrrolidinium and -piperidinium cations, respectively), and Me<sub>4</sub>NTFSI. This implies that the C<sub>2</sub> ↔ C<sub>2</sub> anion disordering mechanism reported here is the norm for plastic crystalline phases of salts with TFSI<sup>-</sup> anions.

**Acknowledgment.** W.A.H. is indebted to the National Science Foundation for the award of a fellowship (International Research Fellowship Program 0202620), and P.C.T. gratefully acknowledges the financial support of the Naval Academy Research Council. Portions of this work were funded by the U.S. Air Force (Air Force Office of Scientific Research). Any opinions, findings, and conclusions or recommendations expressed in this material are those of the author(s) and do not necessarily reflect the views of the National Science Foundation or U.S. Air Force.

**Supporting Information Available:** Experimental procedures (PDF) and X-ray crystallographic data file for Et<sub>4</sub>NTFSI (CIF). This material is available free of charge via the Internet at <http://pubs.acs.org>.

IC0513742

(12) Henderson, W. A.; Young, V. G., Jr.; Passerini, S.; Trulove, P. C.; De Long, H. C. *Chem. Mater.* **2005**, in press.